

Celestial mechanics in Boson Star spacetime

New type of orbits

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Outline

1 Motivations

- GRAVITY
- Tests of General Relativity

2 Boson Star Model

- Field equations
- Numerical solutions
- Astrophysical reality

3 Timelike geodesics in Mini Boson Star model

- Effective potential
- Zero angular momentum stellar orbits

New observations with GRAVITY

GRAVITY instrument

- Optical interferometry in the near-infrared
- astrometric precision of $10 \mu\text{as}$ on each orbit
- Possibility to observe stellar orbits near the Galactic center



Figure: Four 8 m telescopes at VLT (Chile)

Sgr A* : Kerr Black Hole versus Boson Star

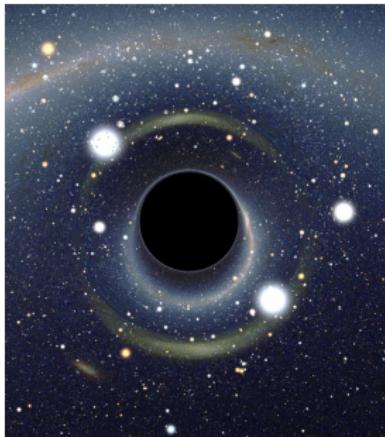


Figure: Image of a Schwarzschild Black Hole

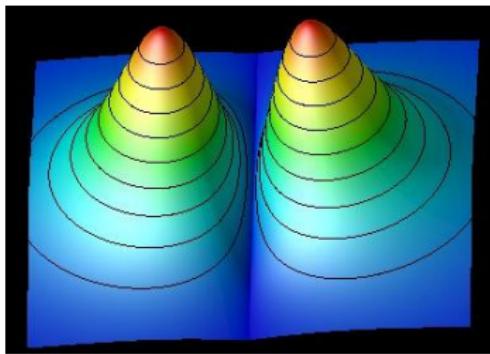


Figure: Rotating Boson Star

Idea : compare the timelike geodesics in those two spacetimes

Field equations

Boson Star : gravitationally bound state of a complex scalar field ϕ which is solution of the following system

- Einstein equations

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = 8\pi T_{\mu\nu}$$

$$T_{\mu\nu} = \frac{1}{2} [\nabla_\mu \bar{\phi} \nabla_\nu \phi + \nabla_\mu \phi \nabla_\nu \bar{\phi}] - \frac{1}{2} g_{\mu\nu} [g^{\gamma\delta} \nabla_\gamma \bar{\phi} \nabla_\delta \phi + V(|\phi|^2)]$$

- Klein Gordon equation

$$\nabla_\mu \nabla^\mu \phi = \frac{dV}{d|\phi|^2} \phi$$

Here we consider “mini” boson stars with $V(|\phi|^2) = \frac{m^2}{h^2} |\phi|^2$

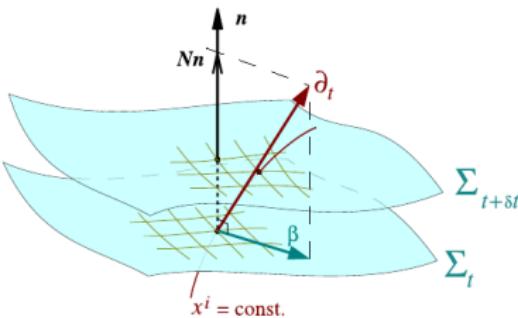
Stationary and axisymmetric solution

Assumptions

- Stationarity and axisymmetry for the spacetime metric $g_{\alpha\beta}$
- Ansatz for the field ϕ

$$\phi = \phi_0(r, \theta) e^{i(\omega t - k\varphi)}$$

with $\phi_0(r, \theta)$ a real function, $\omega \in \mathbb{R}$ and k is an integer.



Solutions found by Kadath using the 3+1 formalism

$$g_{\alpha\beta} dx^\alpha dx^\beta = -N^2 dt^2 + A^2 (dr^2 + r^2 d\theta^2) + B^2 r^2 \sin^2 \theta (d\varphi + \beta^\varphi dt)^2$$

Plots of the boson star field

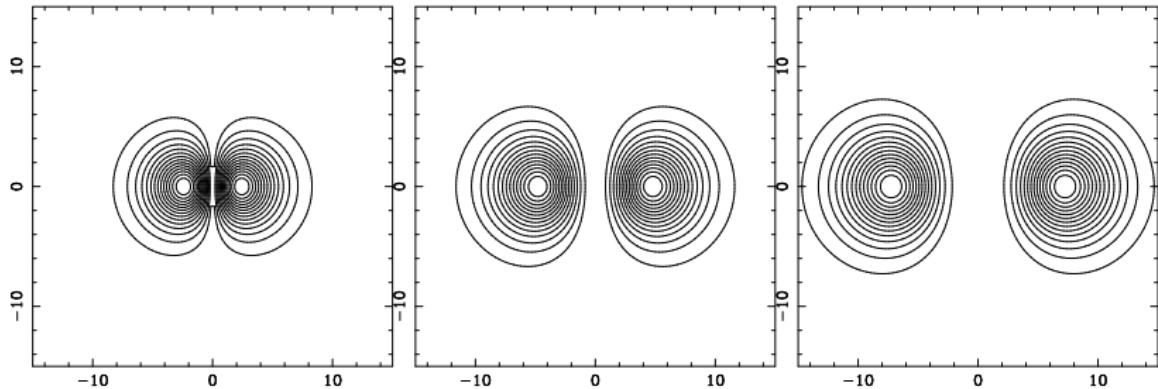


Figure: Isocontours of $\phi_0(r, \theta)$ in the plane $\varphi = 0$ for $\omega = 0.8 m/\hbar$:

$$\phi = \phi_0(r, \theta) e^{i(\omega t - k\varphi)}$$

with $k = 1 ; k = 2 ; k = 3$

Could Sgr A* be a Higgs Star ?

Mass of the Higgs boson

$$m_H = 125.3 \pm 0.6 \text{ GeV}$$

- Mini-boson star $V(|\phi|^2) = \frac{m^2}{h^2} |\phi|^2$

$$M_{crit} = 3 \cdot 10^9 \text{ kg} \ll M_{SgrA*} = 9 \cdot 10^{36} \text{ kg}$$

- Boson star $V(|\phi|^2) = \frac{m^2}{h^2} |\phi|^2 (1 + 2\pi\Lambda |\phi|^2)$ with $\Lambda = 200$:

$$M_{crit} = 8 \cdot 10^{26} \text{ kg} \ll M_{SgrA*}$$

- Solitonic boson star $V(|\phi|^2) = \frac{m^2}{h^2} |\phi|^2 \left(1 - \frac{|\phi|^2}{\sigma^2}\right)^2$ with $\sigma = m_H$:

$$M_{crit} = 4 \cdot 10^{41} \text{ kg} \sim M_{SgrA*}$$

Effective potential

Equation for r :

$$\left(\frac{dr}{d\tau}\right)^2 = \mathcal{V}_{\text{eff}}(r, \epsilon, \ell)$$

with

$$\mathcal{V}_{\text{eff}}(r, \epsilon, \ell) = \frac{1}{A^2} \left[\frac{1}{N^2} (\epsilon + \beta^\varphi \ell)^2 - \frac{\ell^2}{B^2 r^2} - 1 \right]$$

$$\mathcal{V}_{\text{eff}} \geq 0 \Rightarrow \epsilon \geq \epsilon_{\min}$$

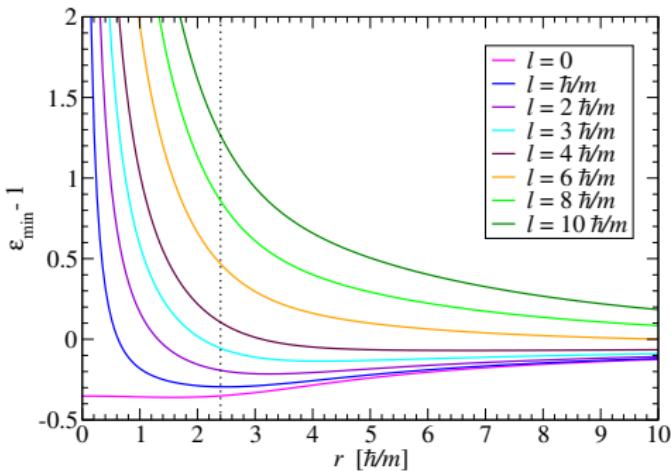
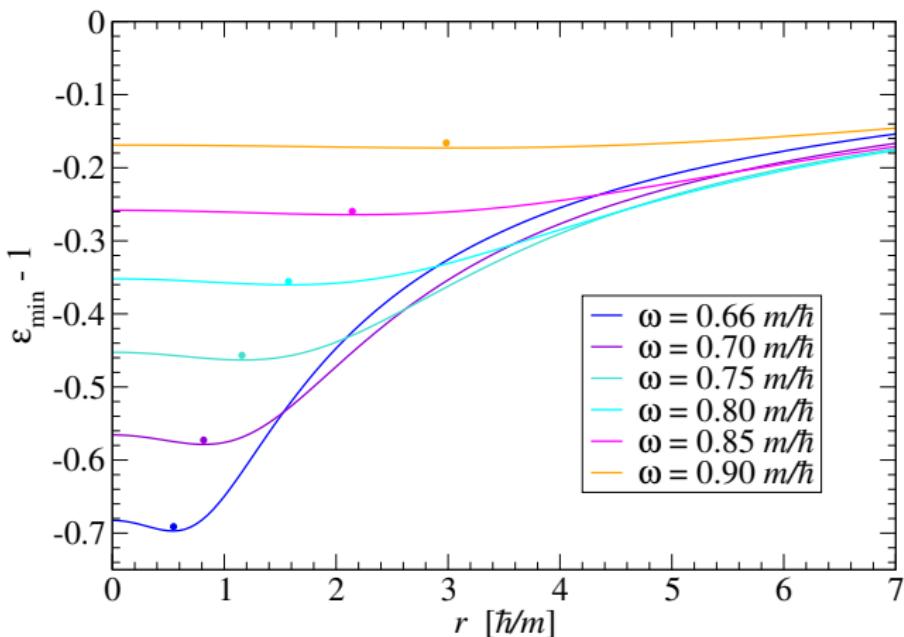


Figure: Effective potential profiles for $\omega = 0.8 \text{ m}/\hbar$ and $k = 1$

Zero angular momentum orbits

Figure: Effective potential for $\ell = 0$ for $k = 1$ and different values of ω

Pointy Petal orbits

Using **GYOTO** ray-tracing code :

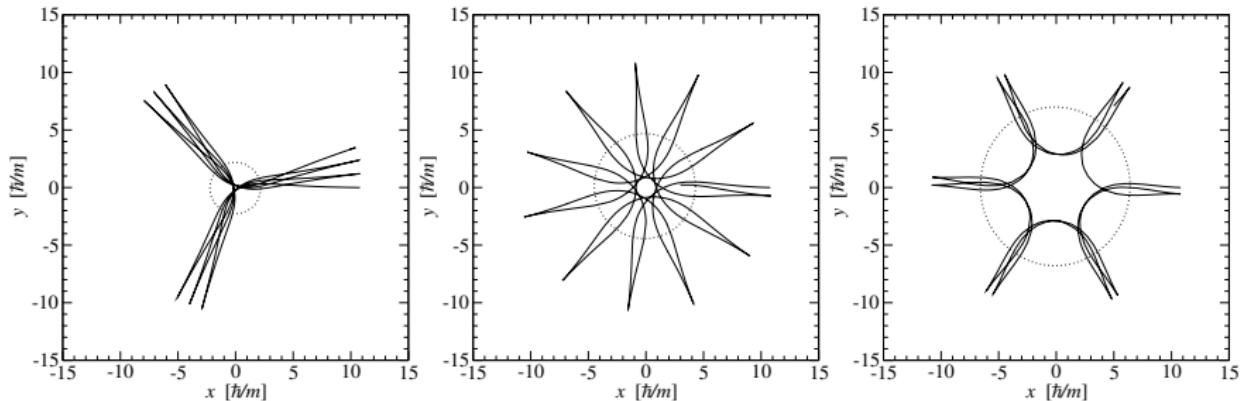


Figure: Orbit of a $\ell = 0$ test particle in the equatorial plane of a boson star with $\omega = 0.8 \text{ m}/\hbar$ and $k = 1, 2, 3$

Pointy Petal orbits

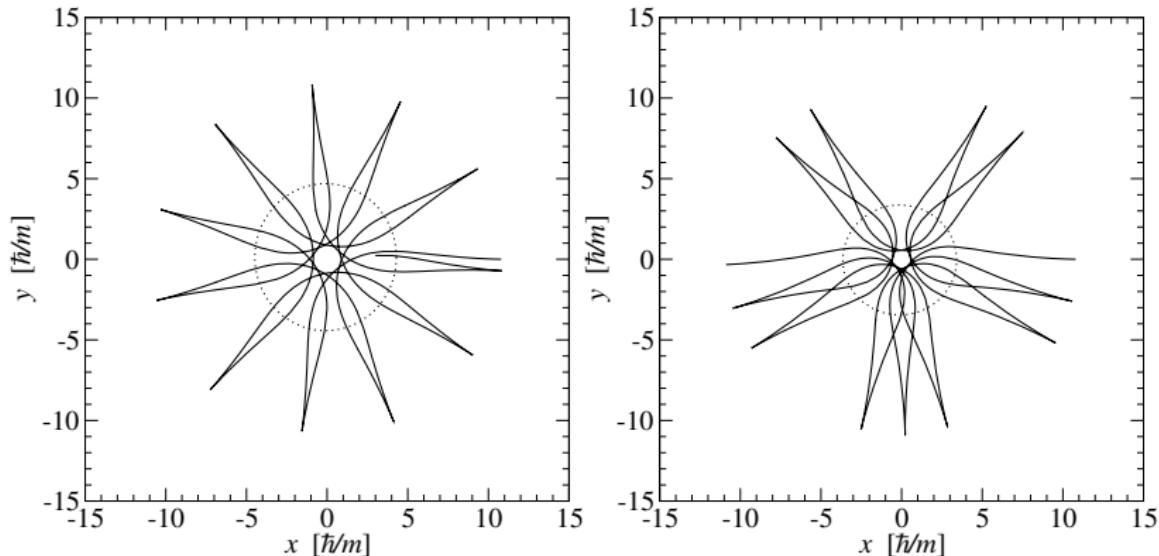


Figure: Orbit of a $\ell = 0$ test particle in the equatorial plane of a boson star with $k = 2$ and $\omega = 0.8 \text{ m}/\hbar$ and $0.75 \text{ m}/\hbar$

Conclusion

New type of orbits with pointy petals

- In Kerr spacetime, orbits with $\ell = 0$ fall into the Black Hole
- If we observe this type of orbits with GRAVITY...
the Galactic Center is definitely a Boson Star !

