

# Black holes and Horndeski's theory

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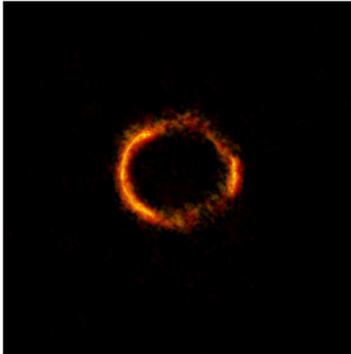
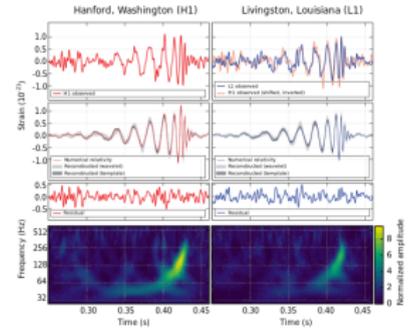
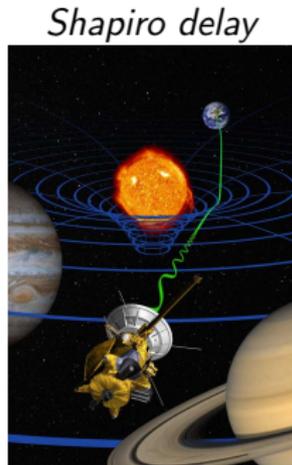
Laboratoire de l'Univers et de ses Théories

GPHYS WORKSHOP  
8/06/18

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## Context

- General Relativity: best, “simplest” classical theory of gravity so far

*Light deflection**Gravitational waves*

...

# Context

- But expected to break down above some energy scale

→ Tested against data from strong gravitational fields:

## LIGO/VIRGO/LISA



*GW from coalescing compact objects*

## EVENT HORIZON TELESCOPE



*Pictures of the surroundings of Sgr A\**

## GRAVITY



*High precision astrometry of the bodies orbiting Sgr A\**

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- Modified gravity: modify GR to account for observed deviations

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- Stationarity of  $S_{GR} \rightarrow G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi T_{\mu\nu}$

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## Horndeski's theory

- $S_{\mathcal{H}} [g, \phi, m^a] = \int \left[ \mathcal{L}_{nonminim}^{(g, \phi)} + \mathcal{L}_{matter} (m^a, \nabla m^a) \right] \sqrt{|\det g|} d^4x$
- Stationarity of  $S_{\mathcal{H}}$   
 $\rightarrow$  most general 2<sup>nd</sup> order Euler-Lagrange equations in  $g$  and  $\phi$

# Cubic Galileon

## Lagrangian formulation of GR

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## Cubic Galileon

$$\mathcal{L}_{nonminim}^{(g, \phi)} = \zeta \left( R^{(g)} - 2\Lambda \right) + (-\eta + \gamma \square \phi) \nabla_{\mu} \phi \nabla^{\mu} \phi$$

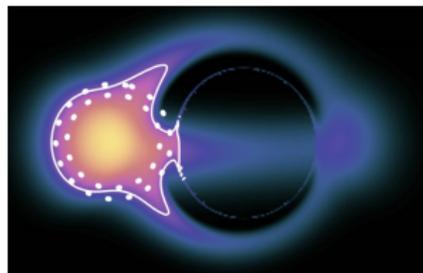
# Cubic Galileon

- Emerges as a limit of an important brane model (“DGP” model)
- Consistent with GW170817:

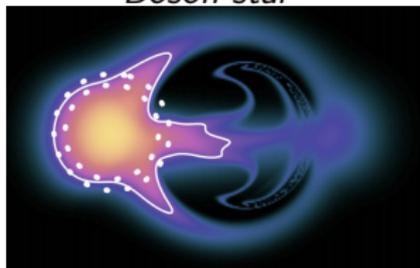
$c_g = c$	$c_g \neq c$
<p>General Relativity quintessence/k-essence Brans-Dicke/<math>f(R)</math> Kinetic Gravity Braiding</p> <p>Viable after GW170817</p>	<p>quartic/quintic Galileons Fab Four de Sitter Horndeski <math>G_{\mu\nu}\phi^\mu\phi^\nu</math>, Gauss-Bonnet</p> <p>Non-viable after GW170817</p>

Ezquiaga & Zumalacárregui, Phys. Rev. Lett. 119, 251304 (2017)

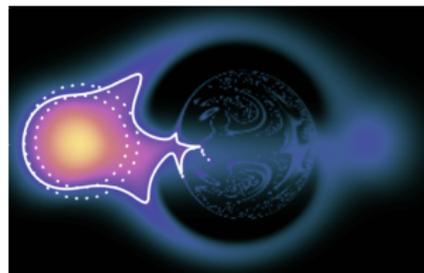
## Predictions of GR



*Kerr black hole*



*Boson star<sup>1</sup>*

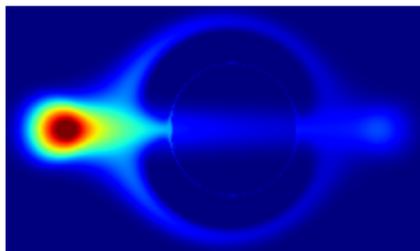


*Black hole  
with scalar hair<sup>2</sup>*

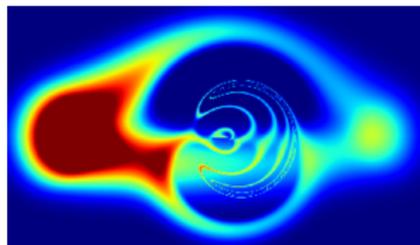
<sup>1</sup>Vincent, Meliani, Grandclément, Gourgoulhon & Straub, *Class. Quantum Grav.* 33, 105015 (2016)

<sup>2</sup>Vincent, Gourgoulhon, Herdeiro & Radu, *Phys. Rev. D* 94, 084045 (2016)  
using the libraries LORENE and KADATH and the ray-tracing code GYOTO

# Predictions of GR



*Regular black hole<sup>3</sup>*



*Rotating naked wormhole<sup>3</sup>*

- *Sgr A\** most likely a black hole

→ Focus on black holes in Cubic Galileon theory

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<sup>3</sup>Lamy, Gourgoulhon, Paumard & Vincent, *Class. Quantum Grav.* 35, 115009 (2018)

## No scalar hair theorems

- Different theory  $\Rightarrow$  different black holes:

e.g. for Cubic Galileon, static spherically symmetric BH with  $\phi(r)$   
 $\implies$  Schwarzschild

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<sup>4</sup>Babichev, Charmousis, Lehébel & Moskalaets, JCAP09(2016)011

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$\longrightarrow$  Introduce a linear time dependence<sup>4</sup>:  $\phi = qt + \Psi(r)$

$\rightarrow$  Preserves spacetime symmetries

$\rightarrow$  Consistent with cosmological dynamics

$\rightarrow$  Yields BH different from GR ones

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<sup>4</sup>Babichev, Charmousis, Lehébel & Moskalets, JCAP09(2016)011

## Additional assumption

- Circularity: stationary axisymmetric spacetime with additional property of orthogonality

$$\iff g_{\mu\nu}(r, \theta) = \begin{pmatrix} -N^2 + B^2\omega^2 r^2 \sin^2 \theta & 0 & 0 & -\omega B^2 r^2 \sin^2 \theta \\ 0 & A^2 & 0 & 0 \\ 0 & 0 & A^2 r^2 & 0 \\ -\omega B^2 r^2 \sin^2 \theta & 0 & 0 & B^2 r^2 \sin^2 \theta \end{pmatrix}$$

→ 4 unknown functions (instead of 10) in 2D

## Problem

- Recall vacuum action of Cubic Galileon:

$$S_{\mathcal{H}}[g, \phi] = \int \left[ \zeta \left( R^{(g)} - 2\Lambda \right) + (-\eta + \gamma \square \phi) \nabla_{\mu} \phi \nabla^{\mu} \phi \right] \sqrt{|\det g|} d^4x$$

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→ Derive the metric equations from  $\frac{\delta S_{\mathcal{H}}}{\delta g_{\mu\nu}}$  + scalar equation from  $\frac{\delta S_{\mathcal{H}}}{\delta \phi}$

- Inject circular metric and scalar ansatz

→ Solve 5 coupled nonlinear PDE's in  $N, A, B, \omega, \Psi$

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→ Solve 5 coupled nonlinear PDE's in  $N, A, B, \omega, \Psi$

- Set boundary conditions defining:
  - event horizon of a rotating BH
  - flat asymptotics

# Spectral methods

- Discretization: consider the truncated decompositions onto standard basis functions

$$\text{e.g. } A(r, \theta) = \sum_{i=0}^{N_r} \sum_{j=0}^{N_\theta} \tilde{A}_{ij} T_i(r) \cos(2j\theta)$$

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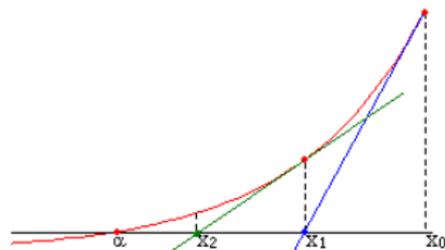
<sup>5</sup>Grandclément, J. Comput. Phys. 229, 3334 (2010), <http://kadath.obspm.fr/>

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- Transforms any system of PDE's into a nonlinear *algebraic* system
- Use Newton-Raphson algorithm implemented in KADATH library<sup>5</sup>



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- Required to be relatively close to the exact solution

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## Initial configuration

- Required to be relatively close to the exact solution
- Build static spherically symmetric solution
  - Non trivial system of ODE's<sup>6</sup>
- Use static BH as initial configuration to obtain a very slowly rotating BH
- Use slowly rotating solution as initial configuration to obtain a less slowly rotating BH...

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<sup>6</sup>Babichev, Charmousis, Lehébel & Moskalets, JCAP09(2016)011

## Summary

- Observations of  $Sgr A^*$  provide new tests of GR
- Horndeski theory such as Cubic Galileon may account for deviations from GR
- A numerical BH would reveal what deviations could be accounted for